

Coherence harvesting with moving qubits

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Our aim

We study the effects of motion on the coherence harvesting protocol [1,4], i.e., the process of extracting quantum superpositions (coherence) from a quantum field and storing it onto another quantum system.

Setup

We consider a two-level system with ground state $|g\rangle$, excited state $|e\rangle$ and energy gap Ω . The qubit is allowed to move along a trajectory $x(\tau) = (t(\tau), x(\tau))$ in 1+1 dimensional flat spacetime—with τ the proper time of the qubit—while interacting with a massless scalar field through the Unruh-DeWitt interaction Hamiltonian [2]

$$\hat{H}_{\text{int}}(\tau) = g\chi(\tau)\hat{\mu}(\tau) \otimes \hat{\phi}(x(\tau)), \quad (1)$$

where g is a small coupling constant, $\chi(\tau)$ is a Gaussian switching function

$$\chi(\tau) = e^{-\frac{\tau^2}{2T^2}}/\sqrt{2\pi T^2}$$

that specifies how the interaction between the qubit and the field is switched on and off (T is the mean interaction duration),

$$\hat{\mu}(\tau) = e^{\Omega\tau}|e\rangle\langle g| + e^{-\Omega\tau}|g\rangle\langle e|$$

is the qubit's monopole moment operator, and

$$\hat{\phi}(x(\tau)) = \int_{-\infty}^{+\infty} \frac{dk}{\sqrt{4\pi|k|}} \left[\hat{a}_k e^{-i(|k|t(\tau) - kx(\tau))} + \text{H.c.} \right]$$

is the field along the qubit's trajectory.

Qubit's motion

We investigate three different states of motion for the qubit: (i) a static qubit with coordinates $t(\tau) = \tau, x(\tau) = 0$, (ii) a qubit moving with constant speed v , following the worldline

$$t(\tau) = \gamma\tau, \quad x(\tau) = \gamma v\tau,$$

where $\gamma = 1/\sqrt{1-v^2}$ is the Lorentz factor, and (iii) a uniformly accelerated qubit, following the hyperbolic trajectory

$$t(\tau) = a^{-1} \sinh(a\tau), \quad x(\tau) = a^{-1} \cosh(a\tau),$$

where a is the acceleration.

Dynamics

We suppose that the qubit starts out in its ground state and the quantum field in a coherent state $|\alpha\rangle$, defined by $\hat{a}_k|\alpha\rangle = \alpha(k)|\alpha\rangle$, where $\alpha(k)$ is a coherent amplitude distribution. We assume a Gaussian profile for the amplitude distribution. The combined qubit-field system evolves from the initial state according to the unitary operator \hat{U} generated by the interaction Hamiltonian (1).

Initial state of the qubit-field system

$$\hat{\rho}_{\text{tot},0} = |g\rangle\langle g| \otimes |\alpha\rangle\langle\alpha|$$

$$\hat{U} = \mathcal{T} \exp \left(-i \int_{-\infty}^{+\infty} \hat{H}_{\text{int}}(\tau) d\tau \right)$$

Final reduced density matrix of the qubit after the interaction

$$\hat{\rho} = \text{tr}_{\hat{\phi}} (U \hat{\rho}_{\text{tot},0} U^\dagger)$$

Quantifying Coherence

To measure the amount of coherence present in the qubit's final state $\hat{\rho}$ we employ the $\ell - 1$ norm of coherence [3], which is equal to the sum of the modulus of its non-diagonal elements

$$C(\rho) = \sum_{i \neq j} |\rho_{ij}|.$$

Coherence Harvesting

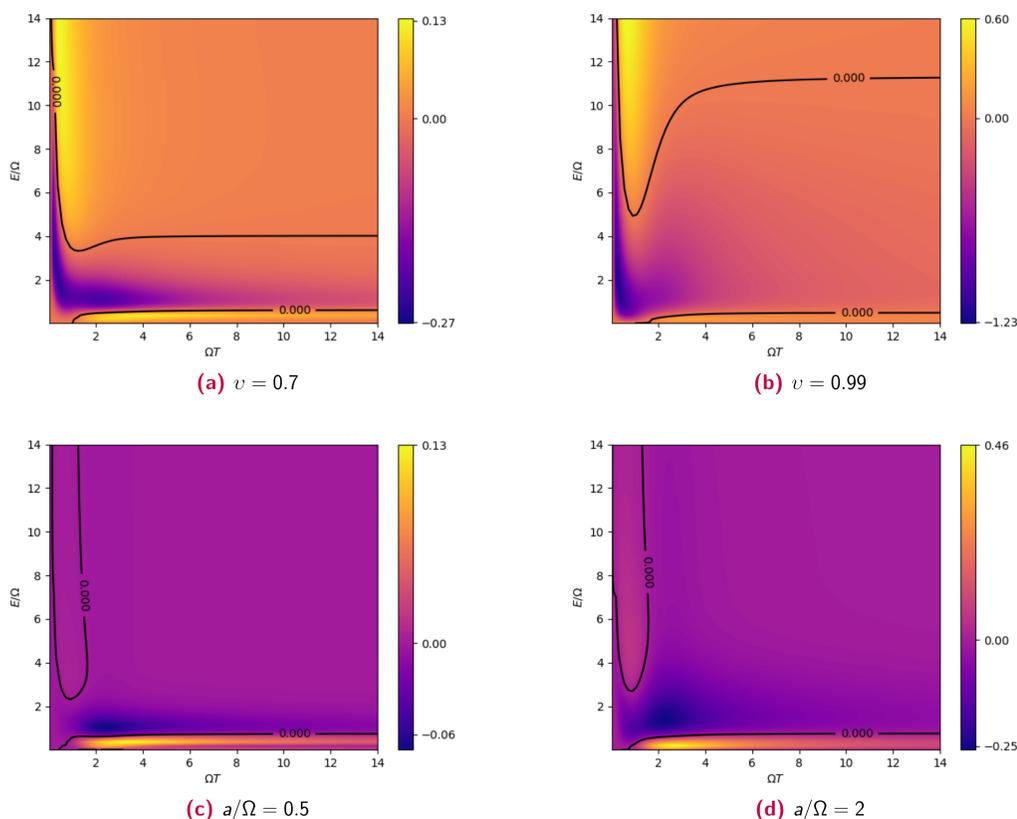


Figure 1: Difference of the amount of coherence extracted scaled by the coupling constant between (a-b) the moving and the static case, $C_v - C_0$, and (c-d) the accelerated and the static case, $C_a - C_0$, as a function of the field's energy E/Ω and ΩT , for various values of speed v and acceleration a . We observe “swelling” regions, where a moving qubit has more coherence than a static one, which are centered around short interaction duration $T \lesssim 1/\Omega$ when $E > \Omega$, and longer interaction times $T \gtrsim 1/\Omega$ when $E < \Omega$.

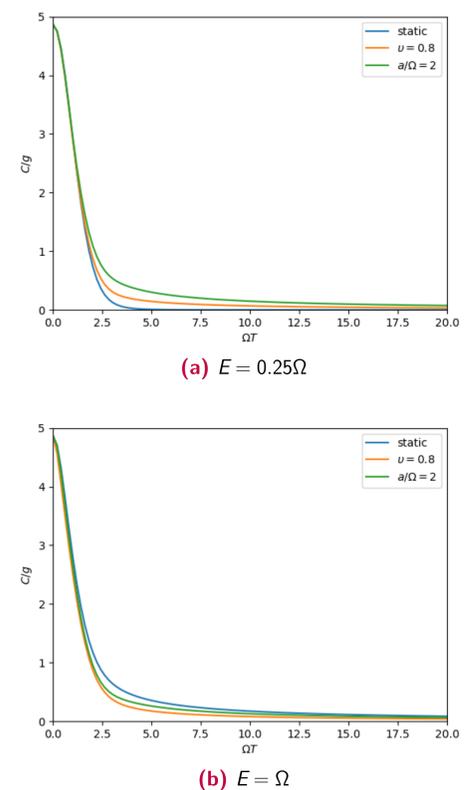


Figure 2: Comparison of the amount of extracted coherence scaled by the coupling constant between a static qubit, a moving qubit with velocity $v = 0.8$, and a uniformly accelerated one with acceleration $a/\Omega = 2$. Depending on the initial energy E of the field, the decoherence rate may be found to be smaller for a moving qubit than a static one.

References

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 [2] N. D. Birrell and P. C. W. Davies, *Quantum Fields in Curved Space*, (Cambridge University Press, 1982). [4] N. K. Kollas and D. Moustos, arXiv:2103.09165.

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